



LAWRENCE
LIVERMORE
NATIONAL
LABORATORY

Laser-Plasma Interactions in High-Energy-Density Plasmas

H. Baldis

October 20, 2006

Disclaimer

This document was prepared as an account of work sponsored by an agency of the United States Government. Neither the United States Government nor the University of California nor any of their employees, makes any warranty, express or implied, or assumes any legal liability or responsibility for the accuracy, completeness, or usefulness of any information, apparatus, product, or process disclosed, or represents that its use would not infringe privately owned rights. Reference herein to any specific commercial product, process, or service by trade name, trademark, manufacturer, or otherwise, does not necessarily constitute or imply its endorsement, recommendation, or favoring by the United States Government or the University of California. The views and opinions of authors expressed herein do not necessarily state or reflect those of the United States Government or the University of California, and shall not be used for advertising or product endorsement purposes.

This work was performed under the auspices of the U.S. Department of Energy by University of California, Lawrence Livermore National Laboratory under Contract W-7405-Eng-48.

Laser-Plasma Interactions in High-Energy-Density Plasmas

Principal Investigator: H. Baldis (University of California, Davis)

High temperature hohlraums (HTH) are designed to reach high radiation temperatures by coupling a maximum amount of laser energy into a small target in a short time. These 400-800 μm diameter gold cylinders rapidly fill with hot plasma during irradiation with multiple beams in 1ns laser pulses. The high-Z plasmas are dense, (electron density, $n_e/n_c \sim 0.1-0.4$), hot (electron temperature, $T_e \sim 10\text{keV}$) and are bathed in a high-temperature radiation field (radiation temperature, $T_{\text{rad}} \sim 300\text{eV}$). Here n_c , the critical density, equals $9 \times 10^{21}/\text{cm}^3$. The laser beams heating this plasma are intense ($\sim 10^{15} - 10^{17} \text{ W/cm}^2$). The coupling of the laser to the plasma is a rich regime for Laser-Plasma Interaction (LPI) physics. The LPI mechanisms in this study include beam deflection and forward scattering. In order to understand the LPI mechanisms, the plasma parameters must be known. An L-band spectrometer is used to measure the electron temperature. A ride-along experiment is to develop the x-radiation emitted by the thin back wall of the half-hohlraum into a thermal radiation source.

Figure 1a shows the experimental setup. About twenty laser beams in three cone angles are incident into a 600 μm diameter, 660 μm long half-hohlraum. The side walls of the hohlraum are gold, usually 20 μm thick. The back wall is thin, $\sim 1 \mu\text{m}$ gold or 1 μm gold overcoated with 1 μm parylene. The high and intermediate angle beams are focused at the center of the Laser Entrance Hole (LEH), but the low angle beams are focused $\sim 250-400 \mu\text{m}$ in front

of the LEH to avoid hitting the back wall. An LPI probe beams is incident almost normal to the hohlraum axis, and aimed to an interaction region, which is the plasma that is 200 μm in front of the LEH. The transmission and forward scatter of this beam is measured with the temporally and spectrally resolved spectrometers and calorimeters in the Full Aperture Backscatter (FABS) diagnostic. Because of the laser beam configuration on the Omega laser, one can use FABS to measure the forward scattered light from opposing beams. If the beam is deflected, it falls onto the NBI plate. A time-averaged image of this deflection is recorded by the NBI camera. The Lband spectrometer views the plasma in the LEH region. The x-radiation emitted by the thin back wall can be used to heat a physics target. To characterize this source, the heating of a witness place placed $\sim 400 \mu\text{m}$ outside the back wall (Figure 1b) was measured.

Beam deflection is measured with the NBI plate. Figure 2 shows images of the NBI plate as a function of LPI probe beam intensity, for two independent interaction beams. As the intensity increases, the beam deflection increases (the cross marks the center of the beam). The LPI beam is “bent” by the plasma flowing out of the target. Beam deflection occurs when the ponderomotively induced density depressions in the plasma move downstream, and carry the light refracted into them. The images from NBI 25 and NBI 30 corresponds to interaction beams B46 and B61, traversing the plasma at angles 31 and 9 degrees respectively, with respect to the normal to the axis of symmetry of the hohlraum. The beam deflections at $5 \times 10^{15} \text{ W/cm}^2$ is approximately 15 and 7.2

degrees respectively. This is the first observation of beam deflection as a function of laser intensity, for different optical paths along the plasma.

Understanding the measured LPI mechanisms depends on knowing the plasma parameters. Radiation-hydrodynamics codes are used to predict the plasma conditions. These must be benchmarked by measurements of n_e and T_e . In highly charged gold, the 3d->2p transitions of individual ionization states are separated by about 40 eV. If these lines can be resolved, the spectrum gives the distribution of the ionization states of gold. This, combined with models that predict the ionization state as a function of electron temperature, would give T_e .

The Lband spectrometer is designed to measure the 3d->2p transitions in gold with high resolution. It is a transmission crystal spectrometer mounted to a single-strip framing camera. It captures a single time and space resolved high resolution spectra. Figure 3a shows a measured spectrum. There are a lump of lines, peaking at 10100 eV, with half-width of about 250 eV. Simulated spectra from the non-local thermodynamic equilibrium (NLTE) code FLYCHK (for $\langle Z \rangle$ as a function of electron temperature) and FLYSPEC (for spectral lines) show similar features: a lump of lines about 200 eV wide. The centroid moves to higher x-ray energy with higher electron temperature. A comparison of the data with simulation shows the measured electron temperature is ~ 7-8 keV.

Measured stimulated Brillouin forward scattering (SBFS) is shown in Figure 4. The SBFS confirms the time at which the plasma reached the interaction region, by the transition from 3w laser light to SBFS. The absence of 3w light after 1ns, may indicate that the nonlinear beam deflection have shifted the beam towards the NBI plate, with the light missing the collecting lens. It is possible that the SBFS is NOT deflected, because of its lower intensity.

The use of the back wall as a radiation source is demonstrated by using it to heat a witness plate. The arrangement for the HTH half-hohlraum with the witness plate is shown in Figure 5. Figure 5a shows a schematic of the target. The witness plate is mounted $\sim 400 \mu\text{m}$ from the back wall, at an 11° tilt to the back wall so that the imaging diagnostic views the witness plate edge-on. The witness plate is a thin chromium foil sandwiched in plastic ($1 \mu\text{m}$ on each side). Figure 5b shows the heated witness plate at two different times. These were taken on two different shots, so the distance of witness plate from back wall is a little different. However, the witness plate is clearly expanding as it is heated. Figure 5c shows the measured width of the witness plate as a function of time from several shots on several different days. The line is the predicted expansion from a rad-hydro simulation, assuming the witness plate is heated solely by radiation from the back wall. The data is consistent with the simulation, which predicts the radiation temperature in the plate to be $\sim 165 \text{ eV}$.

Other measurements of the plasma parameters in the LPI region have been performed. These are (a) the M-shell spectra of gold; (b) raman backscatter and (b) 2w and 4w Thomson scattering. These data are still being

analyzed.

Figure captions:

Figure 1: Experimental set-up: (a) Hot hohlraum showing LPI region, drive, and interaction beams. (b) Hot hohlraum showing location of witness plate and view of Lband spectrometer.

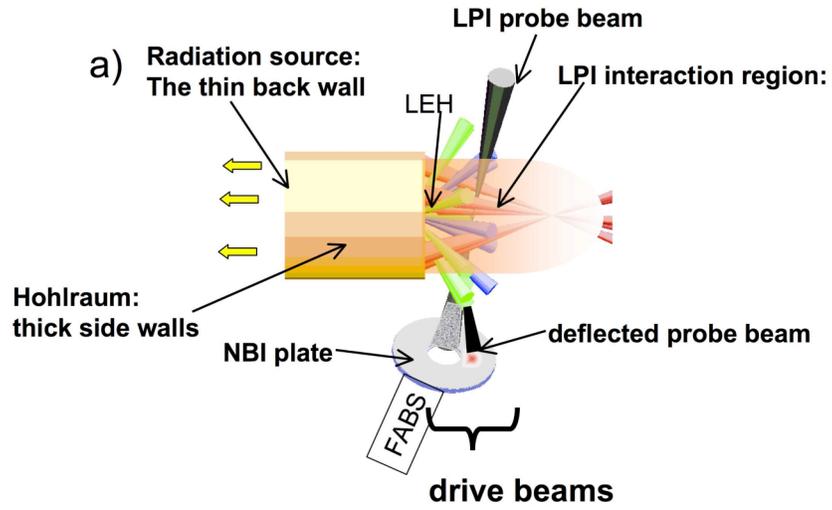
Figure 2: Images of NBI plates show beam deflection as function of LPI probe beam intensity.

Figure 3: 3d-> 3p transitions in highly-charge gold.

Figure 4: Streak camera image showing forward SBS.

Figure 5: Fig 5: New radiation source: thin back wall of HTH half-hohlraum is used to heat a witness plate. (a) The sketch of target shows witness plate mounted $\sim 300\mu\text{m}$ from back wall at 11° angle so it is viewed edge-on by OMEGA diagnostic. (b) The data shows witness plate glowing after it has been heated by back wall (also glowing). (c) measured expansion.

FIGURE 1a



Figure

1b

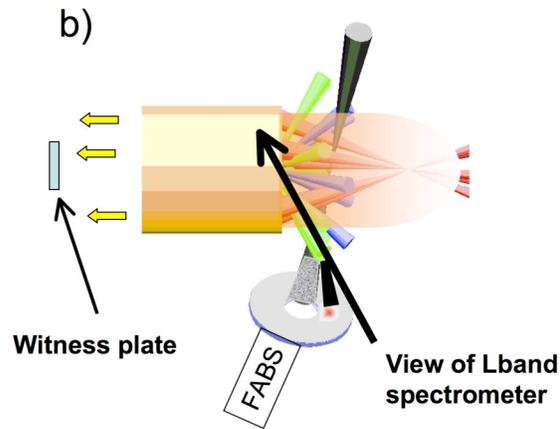


Figure 2

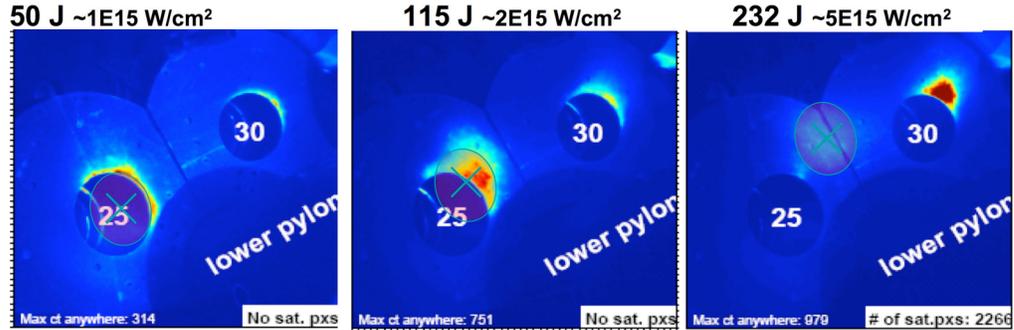
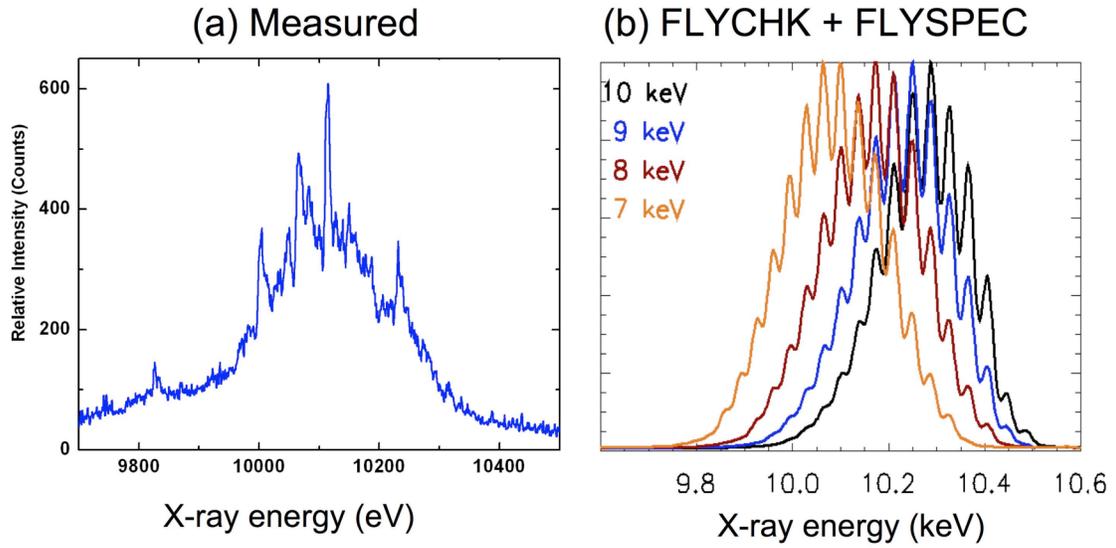


Figure 2: LPI: beam deflection as function of LPI probe beam intensity.

Figure 3



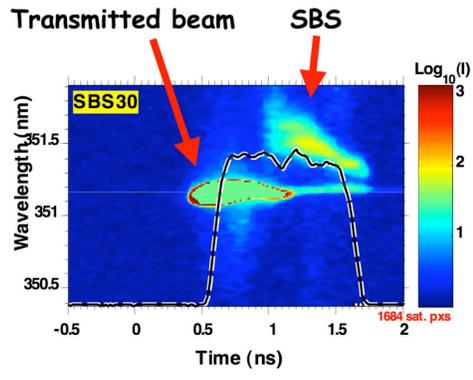


Figure 4

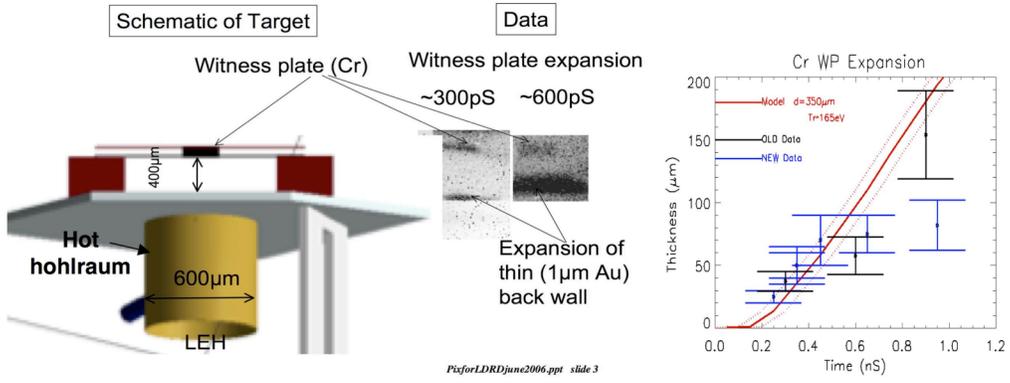


Figure 5